# Extracting Euler Angles from a Rotation Matrix 

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This article attempts to fix a problem which came up when implementing Ken Shoemake's Euler angle extraction in the context of a single-precision floating point library. The original Shoemake code uses double precision, which presumably maintains sufficient precision for the problem not to arise.

We'll follow the notational conventions of Shoemake's "Euler Angle Conversion", Graphics Gems IV, pp. 222-9, with the exception that our vectors are row vectors instead of column vectors. Thus, all our matrices are transposed relative to Shoemake's, and a sequence of rotations will be written from left to right. We'll simplify the discussion by ignoring the various possible axis permutations, and will instead focus on one particular order of applying rotations to illustrate the problem.

Consider the following sequence of rotations: $\theta_{1}$ about the x-axis, then $\theta_{2}$ about the $y$-axis, then $\theta_{3}$ about the z-axis, each rotation being applied about one of the world axes as opposed to one of the body axes. This can be written

$$
\begin{gathered}
R_{x}\left(\theta_{1}\right) R_{y}\left(\theta_{2}\right) R_{z}\left(\theta_{3}\right)=\left(\begin{array}{ccc}
1 & 0 & 0 \\
0 & c_{1} & s_{1} \\
0 & -s_{1} & c_{1}
\end{array}\right)\left(\begin{array}{ccc}
c_{2} & 0 & -s_{2} \\
0 & 1 & 0 \\
s_{2} & 0 & c_{2}
\end{array}\right)\left(\begin{array}{ccc}
c_{3} & s_{3} & 0 \\
-s_{3} & c_{3} & 0 \\
0 & 0 & 1
\end{array}\right) \\
=\left(\begin{array}{ccc}
c_{2} c_{3} & c_{2} s_{3} & -s_{2} \\
s_{1} s_{2} c_{3}-c_{1} s_{3} & s_{1} s_{2} s_{3}+c_{1} c_{3} & s_{1} c_{2} \\
c_{1} s_{2} c_{3}+s_{1} s_{3} & c_{1} s_{2} s_{3}-s_{1} c_{3} & c_{1} c_{2}
\end{array}\right)
\end{gathered}
$$

with $c_{1}=\cos \theta_{1}, s_{1}=\sin \theta_{1}$, etc.

Now suppose we are given a matrix

$$
M=\left(\begin{array}{lll}
m_{00} & m_{01} & m_{02} \\
m_{10} & m_{11} & m_{12} \\
m_{20} & m_{21} & m_{22}
\end{array}\right)
$$

and are required to extract Euler angles corresponding to the above rotation sequence, i.e. find angles $\theta_{1}, \theta_{2}, \theta_{3}$ which make the two matrices equal.

In the general case, Shoemake's code proceeds as follows. First $\theta_{1}$ is extracted using

$$
\begin{gathered}
\theta_{1}=\operatorname{atan} 2\left(m_{12}, m_{22}\right) \\
=\operatorname{atan} 2\left(s_{1} c_{2}, c_{1} c_{2}\right)
\end{gathered}
$$

Next, $c_{2}$ is computed using using

$$
\begin{aligned}
& c_{2}=\sqrt{m_{00}{ }^{2}+m_{01}{ }^{2}} \\
& =\sqrt{c_{2}{ }^{2} c_{3}{ }^{2}+c_{2}{ }^{2} s_{3}{ }^{2}}
\end{aligned}
$$

and hence $\theta_{2}$ using

$$
\theta_{2}=\operatorname{atan} 2\left(-m_{02}, c_{2}\right)
$$

Finally $\theta_{3}$ is obtained using

$$
\begin{gathered}
\theta_{3}=\operatorname{atan} 2\left(m_{01}, m_{00}\right) \\
=\operatorname{atan} 2\left(c_{2} s_{3}, c_{2} c_{3}\right)
\end{gathered}
$$

which is problematic when $m_{00}$ and $m_{01}$ are both very small or zero. We may note that this will also cause $m_{12}$ and $m_{22}$ to be very small or zero (since $m_{02}$ will be near $\pm 1$ ), making the extraction of $\theta_{1}$ equally problematic. Shoemake's solution is to test the value computed for $c_{2}$ against a tiny threshold if it falls below this threshold then the matrix elements reduce to approximately the following:

$$
\left(\begin{array}{ccc}
0 & 0 & -( \pm 1) \\
\pm s_{1} c_{3}-c_{1} s_{3} & \pm s_{1} s_{3}+c_{1} c_{3} & 0 \\
\pm c_{1} c_{3}+s_{1} s_{3} & \pm c_{1} s_{3}-s_{1} c_{3} & 0
\end{array}\right)
$$

and in this case the angles $\theta_{1}$ and $\theta_{3}$ are extracted by a different code path. This matrix provides an example of the phenomenon of gimbal lock, in which the $1^{\text {st }}$ and $3^{\text {rd }}$ axes are brought into alignment by the $2^{\text {nd }}$ rotation, effectively losing a degree of freedom because now $\theta_{1}$ and $\theta_{3}$ act in combination as though they were a single parameter. Shoemake handles this case by forcing $\theta_{3}$ to zero - which is ok, because gimbal lock allows us to arbitrarily set one of $\theta_{1}$ and $\theta_{3}$ and then derive the other. This further reduces the matrix to the following form:

$$
\left(\begin{array}{ccc}
0 & 0 & -( \pm 1) \\
\pm s_{1} & c_{1} & 0 \\
\pm c_{1} & -s_{1} & 0
\end{array}\right)
$$

from which $\theta_{1}$ is easily extracted using

$$
\theta_{1}=\operatorname{atan} 2\left(-m_{21}, m_{11}\right)
$$

This works perfectly well for cases which fall within the small threshold, which is 16 *FLT_EPSILON in Shoemake's code. (I wasn't able to tell where this magic value came from - maybe it's just a very longlived fudge factor.) However, it can be dangerous to use in cases which fall just outside the threshold.

When the routine is passed a real-world single-precision matrix whose elements are subject to some typical rounding errors, $\theta_{1}$ can be expected to take on a fairly chaotic set of values. This is not surprising, as we expect near-gimbal-lock orientations to produce wild jumps in the angle values, and any numerical errors can accentuate these jumps. This in itself is not a major issue because of the close
relationship between $\theta_{1}$ and $\theta_{3}$ near the gimbal lock orientations. Any wobble in the value of $\theta_{1}$ can be counteracted by a suitable anti-wobble in the value of $\theta_{3}$.

However, we have a couple of major expectations of the extracted angles. First, they should approximately reproduce the original matrix when passed to the inverse function. Second, although we may not be able to extract an accurate $\theta_{1}$ or $\theta_{3}$ in isolation, we still expect their resultant angle to be correct: the extracted values should not be independent.

In Shoemake's version, in cases which fall outside the threshold, the angles $\theta_{1}$ and $\theta_{3}$, while algebraically dependent, are numerically independent in the sense that rounding errors in one are not compensated for in the other. In single-precision code these rounding errors can be very large, leading to completely erroneous results when the computed angles are used in an attempt to reconstruct the original matrix.

This theory is easily tested by passing to the Shoemake version a gimbal locked matrix with a couple of the zeros slightly jiggled by amounts just larger than $16^{*}$ FLT_EPSILON - i.e. of the order $10^{-6}$, which could easily arise from rounding errors in real-world single-precision matrices. Sure enough, the result of passing the extracted angles into the inverse function is often a completely different matrix. That's a bug, because we'd like it to be insensitive to rounding errors of order $10^{-6}$. We might try substantially increasing the threshold, to make the results less sensitive to rounding error, but the trouble with this approach is that the approximation used in the gimbal lock cases would no longer be valid.

Have we all simply lived with this problem? Perhaps the errors in double precision matrices tend to be sufficiently small that the bad cases never arise, but this reliance doesn't seem like a robust approach, and is certainly not the way to go for a single-precision library.

Fortunately, there seems to be an easy fix: compute the rotation generated by the first and second extracted angles, and work out the rotation needed in the third angle to match the target matrix. This is easily derived by pre-multiplying the target matrix by the transpose of the reconstructed first-and-second-angle matrix. So we need to compute the following:

$$
\begin{aligned}
& M^{\prime}=\left(\left(\begin{array}{ccc}
1 & 0 & 0 \\
0 & c_{1} & s_{1} \\
0 & -s_{1} & c_{1}
\end{array}\right)\left(\begin{array}{ccc}
c_{2} & 0 & -s_{2} \\
0 & 1 & 0 \\
s_{2} & 0 & c_{2}
\end{array}\right)\right)^{T}\left(\begin{array}{lll}
m_{00} & m_{01} & m_{02} \\
m_{10} & m_{11} & m_{12} \\
m_{20} & m_{21} & m_{22}
\end{array}\right) \\
&=\left(\begin{array}{ccc}
c_{2} & 0 & -s_{2} \\
s_{1} s_{2} & c_{0} & s_{1} c_{2} \\
c_{1} s_{2} & -s_{0} & c_{1} c_{2}
\end{array}\right)^{T}\left(\begin{array}{lll}
m_{00} & m_{01} & m_{02} \\
m_{10} & m_{11} & m_{12} \\
m_{20} & m_{21} & m_{22}
\end{array}\right) \\
&=\left(\begin{array}{ccc}
c_{2} & s_{1} s_{2} & c_{1} s_{2} \\
0 & c_{1} & -s_{1} \\
-s_{2} & s_{1} c_{2} & c_{1} c_{2}
\end{array}\right)\left(\begin{array}{lll}
m_{00} & m_{01} & m_{02} \\
m_{10} & m_{11} & m_{12} \\
m_{20} & m_{21} & m_{22}
\end{array}\right)
\end{aligned}
$$

If this product represents a pure rotation about the z-axis, it must be of the following form:

$$
M^{\prime}=\left(\begin{array}{ccc}
c_{3} & s_{3} & 0 \\
-s_{3} & c_{3} & 0 \\
0 & 0 & 1
\end{array}\right)
$$

and extraction of $\theta_{3}$ should be straightforward. In practice we don't even need to compute a full matrix product. Reading across the middle row of $M^{\prime}$, we obtain

$$
c_{1} m_{10}-s_{1} m_{20}=-s_{3} ; \quad c_{1} m_{11}-s_{1} m_{21}=c_{3}
$$

and hence

$$
\theta_{3}=\operatorname{atan} 2\left(s_{1} m_{20}-c_{1} m_{10}, c_{1} m_{11}-s_{1} m_{21}\right)
$$

which requires that we compute the sine and cosine of the value we extracted for $\theta_{1}$. This way, any gimbal lock instability in the value of $\theta_{1}$ is fed back into the extraction process, and will be counteracted in the value computed for $\theta_{3}$. Note that this method of extracting the $3^{\text {rd }}$ angle makes no assumptions about the other two angles, so that it can be applied in the non-gimbal lock cases too, and no conditional branches are needed.

The final calculation is

$$
\begin{gathered}
\theta_{1}=\operatorname{atan} 2\left(m_{12}, m_{22}\right) \\
c_{2}=\sqrt{m_{00}^{2}+m_{01}^{2}} \\
\theta_{2}=\operatorname{atan} 2\left(-m_{02}, c_{2}\right) \\
s_{1}=\sin \left(\theta_{1}\right), c_{1}=\cos \left(\theta_{1}\right) \\
\theta_{3}=\operatorname{atan} 2\left(s_{1} m_{20}-c_{1} m_{10}, c_{1} m_{11}-s_{1} m_{21}\right)
\end{gathered}
$$

